

Deriv

$$\frac{dI}{ds} = \vec{s} \cdot \nabla I = KI_s - KI$$

$$q = \int_{4\pi} I(\vec{s}) \vec{s} d\Omega = \sum w_i I_i \vec{s}_i$$

Quadrature is used in selecting the Directions $\vec{s}_i$ & weight factors w_i

- use a Symmetric Set invariant to rotations of 90°

- Satisfy moments

$$\int_{4\pi} d\Omega = 4\pi$$

$$\int_{4\pi} \vec{s} d\Omega = 0$$

$$\int_{4\pi} \vec{s} \vec{s} d\Omega = \frac{4\pi}{3} \delta$$

- Also, Satisfy $\int_{\vec{n} \cdot \vec{s} < 0} |\vec{n} \cdot \vec{s}| d\Omega = \pi$ for $\vec{n} = \vec{i}, \vec{j}, \vec{k}$

See table.

$$\vec{S} = (\vec{s}_i \cdot \vec{i}) \vec{i} + (\vec{s}_i \cdot \vec{j}) \vec{j} + (\vec{s}_i \cdot \vec{k}) \vec{k} = s_{ix} \vec{i} + s_{iy} \vec{j} + s_{iz} \vec{k}$$

WSGG

Absorption Spectrum is widely Variable, full of lines.

- Very hard to get an accurate representation of radiative intensities using mean absorption coefficients.
- See slides.

WSGG

Break up / Discretize the Spectrum into bands.

Assume Gray properties in each band.

$$\frac{dI_m}{ds} = -K_m I_m + K_m I_{bm}$$

$$\int_{\eta_j} \frac{dI_m}{ds} = - \int_{\eta_j} K_m I_m d\eta + \int_{\eta_j} K_m I_{bm} d\eta$$

$$\text{Let } I_j = \int_{\eta_j} I_m d\eta, \quad K_m \text{ in band } j \text{ (Gray)} = K_j$$

$$\frac{dI_j}{ds} = -K_j I_j + K_j \underbrace{\frac{\int_{\eta_j} I_{bm} d\eta}{I_b}}_{a_j} * I_b$$

$$\frac{dI_j}{ds} = -K_j I_j + K_j a_j I_b$$

a_j is the fraction of black intensity in band j .

$$\sum a_j = 1$$

$$I = \sum I_j$$

• K_f , a_f are functions of temperature, pressure & composition.

• Bordbar's Model

2) Gray Gases

1 Clear Gas

CO_2 , H_2O mixtures

$$K_f = \sum_{k=0}^4 d_{f,k} M_r^k \quad ; \quad M_r = \tau_{\text{H}_2\text{O}} / \tau_{\text{CO}_2}$$

$$a_f = \sum_{k=0}^4 b_{f,k} T_r^k \quad ; \quad T_r = T / 1200 \text{ K}$$

$$b_{f,k} = \sum_{i=0}^4 c_{f,k,i} M_r^i$$

$d_{f,k}$, $c_{f,k}$ are correlated model coefficients

• Valid for $0.01 \leq M_r \leq 4$

• for M_r outside these bounds, K_f , a_f are interpolated between pure component values for CO_2 or H_2O and the corresponding bound 0.01 or 4.

RCSLW (Rank-Correlated Spectral Line Weighted Sum of Gray Gases) (3)

Even within a band, the gas is not gray, spectrum varies wildly.

→ WSGG above is more like a curve-fit for a given model formulation than an accurate formulation of the model itself given the gray assumption.

The SLW model is a WSGG model for which the gases are gray in a given band by construction.

Instead of integrating over a sequential band of wavenumbers η , SLW integrates a set of η for which the gas cross section is in a given band, which is gray to within the width of the band.

$$(*) \quad \Delta_j = \{ \eta : \tilde{\zeta}_{j-1} < \eta < \tilde{\zeta}_j \} ; \quad \tilde{\zeta}_j \text{ Denotes a Cross-Section boundary between two bands/gases}$$

$$\cdot \text{ Recall } K_\eta = C_\eta N$$

$$\rightarrow \frac{dT_j}{ds} = -K_j I_j + K_j \underbrace{\int_{\Delta_j} I_{b,m} d\eta}$$

evaluate this term.

$$\cdot \text{ before, in WSGG, this was written as } a_j I_b, \text{ with } a_j = \frac{\int_{\Delta_j} I_{b,m} d\eta}{I_b}$$

$$\cdot \text{ Here, it's similar; } \int_{\Delta_j} I_{b,m} d\eta = a_j I_b, \text{ but the bounds of the integration are "over the set of } \eta, \Delta_j$$

$$\cdot a_j \text{ is computed as}$$

$$a_j = F(\tilde{\zeta}_j, \phi, T) - F(\tilde{\zeta}_{j-1}, \phi, T)$$

$$F(L, \phi, T) = \frac{1}{\sigma T^4} \int_{\{ \eta : \kappa_\eta(\eta, \phi) < L \}} E_b(\eta, T) d\eta$$

F is the ALBDF (absorption line blackbody distribution function)

ϕ is the fraction of black emission at some T arising from ~~cross sections~~ η corresponding to cross-sections below some given cross section c .

- F is like a cumulative PDF.
 - Varies from 0 to 1
 - Is monotonic in c (which allows inversion)
 - Can be computed from Spectral Databases,
 - $F = F(c, \phi, T)$ where ϕ is some gas state; T, P, χ_i
 - The T in ϕ can differ from the T in $F(c, \phi, T)$, w/ the latter being the T appearing directly in $\frac{1}{\sigma T^4} \int_{\{m: L_m(m, \phi) < c\}} E_b(m, T) dm$ and the T in ϕ appearing $\nearrow$, which determines the $\int$ range.
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In computing the K_j a logarithmic (Geometric) interpolation is used

$$K_j = N \sqrt{\bar{c}_{j-1} \bar{c}_j}$$

$$\log K_j / N = \frac{\log \bar{c}_{j-1} + \log \bar{c}_j}{2} = \log(\bar{c}_{j-1} \bar{c}_j)^{1/2}$$

$$\rightarrow K_j = (\bar{c}_{j-1} \bar{c}_j)^{1/2} \cdot N$$

Spatial Effects.

The above formulation assumed spatial uniformity of composition, T , where Δ_f is independent of position.

$$\frac{dI_m}{ds} = -k_m I_m + k_m I_{b,m}$$

$$\int_{\Delta_f} \frac{dI_m}{ds} d\eta = - \int_{\Delta_f} k_m I_m d\eta + \int_{\Delta_f} k_m I_{b,m} d\eta$$

$= \frac{d}{ds} \int_{\Delta_f} I_m d\eta$, that is $\frac{d}{ds}$ commutes w/ $\int$, only if the bounds of integration, $\Delta_f = \{\eta : \tilde{c}_{f-1} < c_\eta < \tilde{c}_f\}$ are independent of s , (Position)

• The set of η for given $\tilde{c}$ will vary because $c(\eta)$ Dep on comp ϕ , which Dep on Position.

$$\Delta_f(s_1) \neq \Delta_f(s_2)$$

• $\rightarrow$ Need to account for the Leibnitz terms

- basically "impossible"

- big errors if you Don't.

• Soln is to reformulate so that $\Delta_f(s_1) = \Delta_f(s_2)$

$\rightarrow$ How? $\rightarrow \Delta_f$ is fixed $\rightarrow \tilde{c}_f$ varies w/ position.

Original method is the "Reference Approach"

Assume "ideal" behavior of the spectrum by Defining a reference state ϕ_{ref} w/ the property:

$$\Delta = \{\eta : c_\eta(\phi_{ref}) < c^{ref}\} = \{\eta : c_\eta(\phi_{loc}) < c^{loc}\}$$

$\rightarrow$ AIRDF at c^{loc} & c^{ref} w/ same T_b are equal.

- See figures

(6)

$$F(C^{loc}, \phi_{ref}, T_{ref}) = F(C^{ref}, \phi_{ref}, T_{ref})$$

→ invert to find C^{loc}

Method.

1. Choose ref. state ϕ_{ref}
2. Choose set $\tilde{C}_j^{ref}$ (log-spaced)
3. find local $\tilde{C}_j$, $\tilde{C}_j^{loc}$ by inverting $F(\tilde{C}_j^{loc}, \phi_{loc}, T_{ref}) = F(\tilde{C}_j^{ref}, \phi_{ref}, T_{ref})$
4. $K_j(s) = N_{loc} C_j^{loc} = N_{loc} \sqrt{\tilde{C}_{j-1}^{loc} \tilde{C}_j^{loc}}$
5. $a_j(s) = a_j^{loc} = F(\underbrace{\tilde{C}_j^{ref}}_{\text{ADF ref approach uses loc here}}, \phi_{ref}, T_b = T_{loc}) - F(\underbrace{\tilde{C}_{j+1}^{ref}}_{\text{ADF ref approach uses loc here}}, \phi_{ref}, T_b = T_{loc})$

ADF ref approach uses loc here.

Rank-Correlated SLW model.

2 positions S_1, S_2 w/ states ϕ_1, ϕ_2

Have Spectra $C_n(\eta, \phi_1), C_n(\eta, \phi_2)$

let $\hat{C}_1 = C_n(\hat{\eta}, \phi_1)$ at some given $\hat{\eta}$

$\hat{C}_2 = C_n(\hat{\eta}, \phi_2)$ " " " "

Different $\hat{C}$'s at same $\hat{\eta}$

$$H_1 = \{ \eta : C_n(\eta, \phi_1) < \hat{C}_1 \}$$

$$H_2 = \{ \eta : C_n(\eta, \phi_2) < \hat{C}_2 \}$$

; if $H_1 = H_2$, Spectra are R.C.

R.C. is less restrictive than
The ideal spectrum assumption

See figure.

$$F(\hat{C}_1, \phi_1, T_r) = \frac{1}{\sigma_{T_r}^4} \int_{H_1} E_b(\eta, T_r) d\eta$$

$$F(\hat{C}_2, \phi_2, T_r) = \frac{1}{\sigma_{T_r}^4} \int_{H_2} E_b(\eta, T_r) d\eta$$

$$H_1 = H_2 \rightarrow F(\hat{C}_1, \phi_1, T_r) = F(\hat{C}_2, \phi_2, T_r)$$

if $\hat{C}_1, \phi_1$ are given, invert to get $\hat{C}_2, \phi_2$

Holds for arbitrary ϕ

So, if some F is given or specified, insert to get $\tilde{C}$ at given ϕ

Method

1. Set a T_r as avg T in domain
2. Set $\tilde{F}_j$ with F_j between. These apply to all spatial locations
3. Insert: $\tilde{C}_j = C(\tilde{F}_j, \phi, T_r)$, $C_j = C(F_j, \phi, T_r)$
4. $K_j = C_j R$
5. $a_j = F(\tilde{C}_j, \phi, T_{loc}) - F(\tilde{C}_{j-1}, \phi, T_{loc})$

For mixtures

$$F(C) = F_{CO_2}(C/x_{CO_2}) F_{H_2O}(C/x_{H_2O}) F_{CO}(C/y_{CO})$$

Examples

- See Jupyter